

# Is Howard's Separability Principle a sufficient condition for Outcome Independence?

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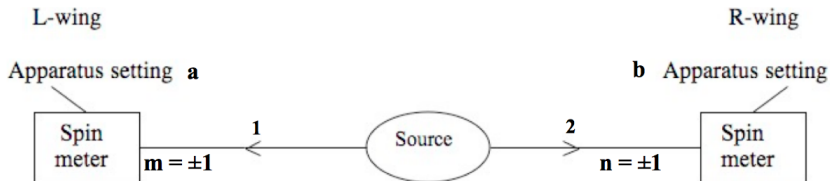
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# Structure

1. Howard's argument
  - 1.1 Set-Up
  - 1.2 Argument
2. Criticisms
  - 2.1 Laudisa
  - 2.2 Fine/Winsberg
  - 2.3 Berkovitz
3. Conclusion

# Setting Up



- $p_{\lambda}^i(m, n | \mathbf{a}, \mathbf{b})$  : Probability of measuring outcome-pair  $(m, n)$  on system(s)  $i$ , given joint state  $\lambda$  and parameter settings  $\mathbf{a}, \mathbf{b}$ .

## Setting Up

All derivations include factorizability [Jarrett, 1984]

$$p_{\lambda}^{12}(m, n|\mathbf{a}, \mathbf{b}) = p_{\lambda}^1(m|\mathbf{a}) \cdot p_{\lambda}^2(n|\mathbf{b})$$

Equivalent to conjunct of:

- Parameter Independence (PI)

$$p_{\lambda}^1(m|\mathbf{a}, \mathbf{b}) = p_{\lambda}^1(m|\mathbf{a})$$

$$p_{\lambda}^2(n|\mathbf{a}, \mathbf{b}) = p_{\lambda}^2(n|\mathbf{a})$$

- Outcome Independence (OI)

$$p_{\lambda}^1(m|\mathbf{a}, \mathbf{b}, n) = p_{\lambda}^1(m|\mathbf{a}, \mathbf{b})$$

$$p_{\lambda}^2(n|\mathbf{a}, \mathbf{b}, m) = p_{\lambda}^2(n|\mathbf{a}, \mathbf{b})$$

How to physically interpret the two?

## The argument

- Einstein articulates two principles that let him regard QM as incomplete: Locality and Separability

***LP:** No supraluminal signalling between space-like separated space-time regions.*

***SP:** The contents of any two regions of space-time separated by a non-vanishing spatio-temporal interval constitute separable physical systems, in the sense that each possesses its own, distinct physical state **[SP(1)]**, and the joint state of the two systems is wholly determined by these separate states **[SP(2)]** [Howard, 1989, 226]*

- Locality  $\leftrightarrow$  PI, Separability  $\leftrightarrow$  OI

## Formalizing separability

### State separability condition [Howard, 1989]

Two systems 1 and 2 are “state separable” if there exist distinct states  $\alpha$  and  $\beta$  for these systems s.t.

$$p_{\lambda}^{12}(m, n|\mathbf{a}, \mathbf{b}) = p_{\alpha}^1(m|\mathbf{a}, \mathbf{b}) \cdot p_{\beta}^2(n|\mathbf{a}, \mathbf{b}) \quad (\text{SEP})$$

**State:** *A conditional probability measure assigning probabilities to outcomes conditional upon the presence of global measurement contexts. ( $\alpha : m \rightarrow p_{\alpha}(m|\mathbf{a}, \mathbf{b}), m \in M$ )*

- Together with identifications

$$p_{\alpha}^1(m|\mathbf{a}, \mathbf{b}) = p_{\lambda}^1(m|\mathbf{a}, \mathbf{b}) \quad (\text{ID I})$$

$$p_{\beta}^2(n|\mathbf{a}, \mathbf{b}) = p_{\lambda}^2(n|\mathbf{a}, \mathbf{b}), \quad (\text{ID II})$$

easily proven that  $SEP \equiv OI \quad \therefore \text{no OI} \rightarrow \text{no SP.}$

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# Laudisa [1995]- Unsoundness of the proof

$$p_{\alpha}^1(m|\mathbf{a},\mathbf{b}) = p_{\lambda}^1(m|\mathbf{a},\mathbf{b}) \quad (\text{ID I})$$

$$p_{\beta}^2(n|\mathbf{a},\mathbf{b}) = p_{\lambda}^2(n|\mathbf{a},\mathbf{b}), \quad (\text{ID II})$$

- Laudisa [1995] argues that Howard's equivalence proof is unsound, because ID I/II not generally true.
- State component distinction [van Fraassen, 1991]
  - *Value* component: “fully specified by stating which observables have values and what they are”
  - *Dynamic* component: “fully specified by stating how the system will develop if isolated, and how if acted upon in any definite, given fashion”
- Measurement is interaction  $\rightarrow$  outcome probabilities comprised in the dynamic component of a state, may differ between  $p_{\alpha}$  and  $p_{\lambda}$ .



# Rebuttal

$$p_{\alpha}^1(m|\mathbf{a},\mathbf{b}) = p_{\lambda}^1(m|\mathbf{a},\mathbf{b}) \quad (\text{ID I})$$

$$p_{\beta}^2(n|\mathbf{a},\mathbf{b}) = p_{\lambda}^2(n|\mathbf{a},\mathbf{b}), \quad (\text{ID II})$$

- However, reference to system via SP  
( $\alpha^i : m \rightarrow p_{\alpha}^i(m|\mathbf{a}, \mathbf{b}), m \in M$ )
- Detectors in different space-time regions  $\rightarrow$  Two *systems* with distinct states  $\alpha^1$  and  $\beta^2$

# Rebuttal

$$p_{\alpha}^1(m|\mathbf{a},\mathbf{b}) = p_{\lambda}^1(m|\mathbf{a},\mathbf{b}) \quad (\text{ID I})$$

$$p_{\beta}^2(n|\mathbf{a},\mathbf{b}) = p_{\lambda}^2(n|\mathbf{a},\mathbf{b}), \quad (\text{ID II})$$

- Formally summarised in ID I/II: Probability for measuring an outcome at some space-time region is completely determined by the state of the system in that region.

$$SP \rightarrow (SEP) \rightarrow \rho_{\lambda}^{12} = \rho_{\alpha}^1 \otimes \rho_{\beta}^2 \xrightarrow{\text{tr}_2(\cdot)} \rho_{\lambda}^1 = \rho_{\alpha}^1$$

$\Rightarrow$  By reference to marginal systems ID I/II entailed by  $(SEP)$ .

## Information “at” margins

“This is because  $\alpha$ , by SP, captures exclusively and exhaustively all the information that  $\lambda$  contains about system 1.”

- “...one should not understand the transmission of information<sub>t</sub> on the model of transporting potatoes.” (Timpson, 2006)
- Probability distributions the same,  $S(\rho_\lambda^1) = S(\rho_\alpha^1)$  if (SEP) true
- No recourse to information concept required for argument

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## Fine and Winsberg [2003]- The minimalist *non sequitur*

- Disagree with Howard's formalization: SP does not specify how to formally determine the joint state.
- Can produce joint states that satisfy the separability principle *and* violate Bell inequalities.

$F(..)$

1. Given a joint state  $\lambda$ , first identifies *bijectively* “marginal distributions”  $p_\lambda(m|\mathbf{a}), p_\lambda(n|\mathbf{b})$  [SP(1)]
2. Finds two-placed function  $F(..)$  that completely reproduces the joint state distribution [SP(2)]

## Two possible counters

- Expand argument to factorizability?
- Restrict admissible  $F(..)$  to include only factorisable  $F(..)$ ?
  - Minimalism: Violation of Bell inequalities is *mathematical* feature of QM due to the inability to define joint probabilities for incompatible observables.
  - Muller and Placek [2001]: Argument against minimalism involving OI- and PI-like assumptions
    - Motivate restriction of mathematical freedom in devising  $F(..)$ , based on *physical* principles.
    - Nothing done to motivate exclusion of non-factorisable  $F(..)$ .

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## Berkovitz [1998]- Well localized particles

- Howard's equivalence proof “presupposes that particles in Bell-type experiments are well localized during the measurements”.
- SP does not imply this  $\rightarrow$  OI not entailed by SP.

### Being well-localized (WL)

A system  $A$  is well localized in space-time iff for any space-time region the question whether  $A$  is constituted by the contents (not necessarily completely) of this region has a well-defined answer and this answer is either “yes” or “no”.



## Counter?

- Is WL required? Yes. “...contents of ...regions of space-time ...constitute physical systems in the sense that each possess ...distinct state.”

⇒ Spatiotemporal criterion of distinctness/individuation.

- Can SP be taken to imply WL? No.
- Can we draw same conclusion? Maybe
  - WL treated as given
  - Question about applicability of SP: Can there be *meaningful* outcome-dependence in situations where SP does not apply?
    - Outcome dependent dynamics collapse (CSL) theories (e.g. GRW)?

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# Conditions for upholding the argument

- Committed to:
  - States being dynamical states as discussed by van Fraassen.
  - Subscribe to feasibility of Müller and Placek's framework (SOBST)  $\Rightarrow$  orthodox QM incomplete
  - Distinctness of states implies their being well localized (spatio-temporal criterion of individuation)
  - Non-feasibility of outcome-dependent dynamic collapse theories (GRW, CSL)
- Strong separability principle  $\rightarrow$  little gain in knowing its falsehood.

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